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# On uniqueness of the complex <br> Monge-Ampère equation on compact Kähler manifolds 

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# ON UNIQUENESS OF THE COMPLEX <br> MONGE-AMPÈRE EQUATION ON COMPACT KÄHLER MANIFOLDS 

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#### Abstract

We prove a partial uniqueness for solutions of the complex Monge-Ampère equation on a compact Kähler manifold in the class of quasiplurisubharmonic functions introduced recently by Guedj and Zeriahi.


## 1. Introduction

Let $(X, \omega)$ be a compact Kähler manifold of complex dimension $n . \operatorname{PSH}(X, \omega)$ will denote the class of quasiplurisubharmonic functions $\varphi$ on $X$ satisfying $\omega_{\varphi}:=$ $\omega+d d^{c} \varphi \geq 0$. Guedj and Zeriahi [9] introduced the class $\mathcal{E}(X, \omega)$ of functions $\varphi \in \operatorname{PSH}(X, \omega)$ satisfying $\int_{\{\varphi>-\infty\}} \omega_{\varphi}^{n}=\int_{X} \omega^{n}$ (the measure $\omega_{\varphi}^{n}=\omega_{\varphi} \wedge \cdots \wedge \omega_{\varphi}$ is well defined on $\{\varphi>-\infty\}$ for any $\varphi \in \operatorname{PSH}(X, \omega)$ by [1]). They showed in particular that for $\varphi \in \mathcal{E}(X, \omega)$ the measure $\omega_{\varphi}^{n}$ is well defined on $X$ (with total mass $\int_{X} \omega^{n}$ ), vanishes on pluripolar sets, and is continuous (in the weak ${ }^{*}$ topology) for decreasing sequences in $\mathcal{E}(X, \omega)$.

One of the main results in [9], building on earlier work of Yau [11], Kołodziej [10], and Cegrell [6] was the existence of a solution of the Dirichlet problem

$$
\left\{\begin{array}{l}
\varphi \in \mathcal{E}(X, \omega)  \tag{1}\\
\omega_{\varphi}^{n}=\mu \\
\max _{X} \varphi=0
\end{array}\right.
$$

provided that $\mu$ is a measure on $X$ vanishing on pluripolar sets and with total mass $\int_{X} \omega^{n}$ (which are of course necessary conditions).

The uniqueness in (1) was posed as a problem in [9]. It had been proved in [3] for bounded $\varphi, \psi$ (for $X=\mathbb{P}^{n}$ it had been earlier done in [2] with much more complicated methods). As observed in [9] the method from [3] actually gives the following result: if $\varphi \in \mathcal{E}(X, \omega)$ and $\psi \in \mathcal{E}^{1}(X, \omega)$ (where $\mathcal{E}^{p}(X, \omega):=\{\psi \in \mathcal{E}(X, \omega)$ : $\left.\left.\int_{X}|\psi|^{p} \omega_{\psi}^{n}<\infty\right\}, p>0\right)$ are such that $\omega_{\varphi}^{n}=\omega_{\psi}^{n}$ then $\varphi-\psi=$ const.

[^0]The goal of this note is to prove the following improvement:
Theorem 1. Assume $\varphi \in \mathcal{E}(X, \omega)$ and $\psi \in \mathcal{E}^{p}(X, \omega)$ for some $p>1-2^{1-n}$. If $\omega_{\varphi}^{n}=\omega_{\psi}^{n}$ then $\varphi-\psi=$ const.

Let us mention that if $\varphi \in \mathcal{E}(X, \omega)$ are such that $\omega_{\varphi}^{n}=\omega_{\psi}^{n}=: \mu$ then we must have $\omega_{\max \{\varphi, \psi\}}^{n}=\mu$ (see [9], Proposition 3.4, we will make use of this result in the proof of Theorem 1) and also $\omega_{t \varphi+(1-t) \psi}^{n}=\mu, 0 \leq t \leq 1$ (see [8]).

In [4] and [5] the class $\mathcal{D}$ of (germs of) plurisubharmonic functions was defined (it was shown that it is actually the same as the class $\mathcal{E}$ studied in [7]). It is the maximal subclass of the class of plurisubharmonic functions where the complex Monge-Ampère operator $\left(d d^{c}\right)^{n}$ can be defined (as a regular measure) so that it is continuous for decreasing sequences. It was shown in [4] that $\mathcal{D}=P S H \cap W_{l o c}^{1,2}$ for $n=2$ and it was characterized (similarly, but in a more complicated way) for $n \geq 3$ in [5].

A natural counterpart of the class $\mathcal{D}$ on a compact Kähler manifold is the class $\mathcal{D}(X, \omega)$ consisting of those $\varphi \in \operatorname{PSH}(X, \omega)$ such that locally $\varphi+g \in \mathcal{D}$, where $g$ is a local potential for $\omega$ (that is $\omega=d d^{c} g$ ). In particular, for $n=2$ we get $\mathcal{D}(X, \omega)=$ $\mathcal{P} S H(X, \omega) \cap W^{1,2}(X)$ (and $\subset$ for arbitrary $n$ ). The measure $\omega_{\varphi}^{n}$ is of course well defined for $\varphi \in \mathcal{D}(X, \omega)$. By $\mathcal{D}_{a}(X, \omega)$ denote the class of those $\varphi \in \mathcal{D}(X, \omega)$ for which $\omega_{\varphi}^{n}$ vanishes on pluripolar sets. It follows that $\mathcal{D}_{a}(X, \omega) \subset \mathcal{E}(X, \omega)$ but by Example 2.14 in [9] we don't have the equality in general.

By Lemma 5.14 in [7] the Dirichlet problem (1), where $\mu$ is a measure on $X$ vanishing on pluripolar sets and with total mass $\int_{X} \omega^{n}$, always has a local solution in $\mathcal{D}_{a}$. This is therefore perhaps natural to ask whether it has a global solution belonging to $\mathcal{D}_{a}(X, \omega)$, which would be an improvement of Theorem A in [9]. However, using Theorem 1 we can show that this is not the case:

Theorem 2. Let $(X, \omega)$ be the projective space $\mathbb{P}^{n}$ with the Fubini-Study metric. There exists a measure $\mu$ on $X$, vanishing on pluripolar sets and with total mass $\int_{X} \omega^{n}$, such that there is no $\varphi \in \mathcal{E}(X, \omega) \cap W^{1,2}(X)$ satisfying $\omega_{\varphi}^{n}=\mu$.

In the proofs of Theorems 1 and 2 we will follow the notation from [9] and use various results proved in that article. We always assume that $(X, \omega)$ is a fixed Kähler manifold.

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## Proofs

As the proof in dimension 2 is simpler and more transparent, we first prove Theorem 1 in this case.

Proof of Theorem 1 in dimension 2. If $\widetilde{\psi}:=\max \{\varphi, \psi\}$ then $\widetilde{\psi} \geq \psi, \widetilde{\psi} \in \mathcal{E}^{p}(X, \omega)$ (by Lemma 2.3 in [9]), and by Proposition 3.4 in [9] we have $\omega_{\tilde{\psi}}^{2}=\omega_{\psi}^{2}=\omega_{\varphi}^{2}$. We may thus assume that $\varphi \leq \psi \leq-1$. Then, if we set $\psi^{j}:=\max \{\varphi, \psi-j\}$, we have $\psi-j \leq \psi^{j} \leq \psi, \psi^{j} \in \mathcal{E}^{p}(X, \omega)$, and $\psi^{j}$ decreases to $\varphi$ as $j \rightarrow \infty$. Without loss of generality we may thus assume that $0 \leq \rho:=\psi-\varphi \leq C$; then both $\varphi$ and $\psi$ belong to $\mathcal{E}^{p}(X, \omega)$.

We now set $\varphi_{j}:=\max \{\varphi,-j\}, \psi_{j}:=\max \{\psi,-j\}, \rho_{j}:=\psi_{j}-\varphi_{j}$, and $h_{j}:=$ $\left(\varphi_{j}+\psi_{j}\right) / 2$. First, we claim that

$$
\begin{equation*}
\lim _{j \rightarrow \infty} \int_{X} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega_{h_{j}}=0 \tag{2}
\end{equation*}
$$

We have

$$
\int_{X} \rho_{j}\left(\omega_{\varphi_{j}}^{2}-\omega_{\psi_{j}}^{2}\right)=-2 \int_{X} \rho_{j} d d^{c} \rho_{j} \wedge \omega_{h_{j}}=2 \int_{X} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega_{h_{j}} .
$$

On the other hand,

$$
\begin{aligned}
\left|\int_{X} \rho_{j}\left(\omega_{\varphi_{j}}^{2}-\omega_{\psi_{j}}^{2}\right)\right| & =\left|\int_{\{\varphi \leq-j\}} \rho_{j}\left(\omega_{\varphi_{j}}^{2}-\omega_{\psi_{j}}^{2}\right)\right| \\
& \leq C\left(\int_{\{\varphi \leq-j\}} \omega_{\varphi_{j}}^{2}+\int_{\{\psi \leq-j\}} \omega_{\psi_{j}}^{2}\right) \rightarrow 0
\end{aligned}
$$

We thus get (2).
Set $\chi(t):=-\sqrt{-t}, t \leq-1$. We want to show the following improvement of (2)

$$
\begin{equation*}
\lim _{j \rightarrow \infty} \int_{X} \chi \circ h_{j} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega_{h_{j}}=0 \tag{3}
\end{equation*}
$$

Similarly as above we have

$$
\begin{align*}
\left|\int_{X} \chi \circ h_{j} \rho_{j}\left(\omega_{\varphi_{j}}^{2}-\omega_{\psi_{j}}^{2}\right)\right| & =\left|\int_{\{\varphi \leq-j\}} \chi \circ h_{j} \rho_{j}\left(\omega_{\varphi_{j}}^{2}-\omega_{\psi_{j}}^{2}\right)\right|  \tag{4}\\
& \leq C|\chi(-j)|\left(\int_{\{\varphi \leq-j\}} \omega_{\varphi_{j}}^{2}+\int_{\{\psi \leq-j\}} \omega_{\psi_{j}}^{2}\right) \rightarrow 0
\end{align*}
$$

because $\varphi, \psi \in \mathcal{E}^{p}(X, \omega)$ and $p>1 / 2$. On the other hand,

$$
\int_{X} \chi \circ h_{j} \rho_{j}\left(\omega_{\varphi_{j}}^{2}-\omega_{\psi_{j}}^{2}\right)=2 \int_{X} d\left(\chi \circ h_{j} \rho_{j}\right) \wedge d^{c} \rho_{j} \wedge \omega_{h_{j}} .
$$

By (4) it is enough to estimate, using the Schwarz inequality,

$$
\begin{aligned}
\mid \int_{X} \rho_{j} \chi^{\prime} \circ h_{j} d h_{j} & \wedge d^{c} \rho_{j} \wedge \omega_{h_{j}} \mid \\
& \leq C \sqrt{\int_{X} \chi^{\prime} \circ h_{j} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega_{h_{j}}} \sqrt{\int_{X} \chi^{\prime} \circ h_{j} d h_{j} \wedge d^{c} h_{j} \wedge \omega_{h_{j}}}
\end{aligned}
$$

In order to show that the last integral is bounded in $j$ we write

$$
\begin{aligned}
\int_{X} \chi^{\prime} \circ h_{j} d h_{j} \wedge d^{c} h_{j} \wedge \omega_{h_{j}} & =-\int_{X} \chi \circ h_{j} d d^{c} h_{j} \wedge \omega_{h_{j}} \\
& \leq-\int_{X} \chi \circ h_{j} \omega_{h_{j}}^{2} \\
& \leq-4 \int_{X} \chi \circ \varphi_{j} \omega_{\varphi_{j}}^{2}
\end{aligned}
$$

by Lemma 2.3 in [9]. Now from (2) (note that $0 \leq \chi^{\prime} \leq 1$ ) we thus get (3).
Proceeding as in [3] we write

$$
\int_{X} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega=\int_{X} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega_{h_{j}}-\int_{X} d \rho_{j} \wedge d^{c} \rho_{j} \wedge d d^{c} h_{j}
$$

so by (2) it is enough to estimate the last integral. We have

$$
-\int_{X} d \rho_{j} \wedge d^{c} \rho_{j} \wedge d d^{c} h_{j}=\int_{X} d \rho_{j} \wedge d^{c} h_{j} \wedge d d^{c} \rho_{j}=\int_{X} d \rho_{j} \wedge d^{c} h_{j} \wedge\left(\omega_{\psi_{j}}-\omega_{\varphi_{j}}\right)
$$

By the Schwarz inequality and since $\omega_{\psi_{j}} \leq 2 \omega_{h_{j}}$

$$
\begin{aligned}
& \left|\int_{X} d \rho_{j} \wedge d^{c} h_{j} \wedge \omega_{\psi_{j}}\right| \\
& \quad \leq 2 \sqrt{\int_{X} \frac{1}{\chi^{\prime} \circ h_{j}} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega_{h_{j}}} \sqrt{\int_{X} \chi^{\prime} \circ h_{j} d h_{j} \wedge d^{c} h_{j} \wedge \omega_{h_{j}}}
\end{aligned}
$$

The last integral is bounded in $j$. In our case we also have $1 / \chi^{\prime}=2 \chi$ and by (3) we get

$$
\lim _{j \rightarrow \infty} \int_{X} d \rho_{j} \wedge d^{c} h_{j} \wedge \omega_{\psi_{j}}=0
$$

Similarly we show that

$$
\lim _{j \rightarrow \infty} \int_{X} d \rho_{j} \wedge d^{c} h_{j} \wedge \omega_{\varphi_{j}}=0
$$

and thus

$$
\lim _{j \rightarrow \infty} \int_{X} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega=0
$$

For the proof of Theorem 1 in arbitrary dimension we will need some preparatory results.
Lemma 3. For $p>0, k=1, \ldots, n$, and $\varphi \in \operatorname{PSH}(X, \omega) \cap L^{\infty}(X)$ with $\varphi \leq-1$ we have

$$
\begin{equation*}
\int_{X}(-\varphi)^{p} \omega_{\varphi}^{n-k} \wedge \omega^{k} \leq \int_{X}(-\varphi)^{p} \omega_{\varphi}^{n} \tag{5}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{X}(-\varphi)^{p-1} d \varphi \wedge d^{c} \varphi \wedge \omega_{\varphi}^{n-k} \wedge \omega^{k-1} \leq \frac{1}{p} \int_{X}(-\varphi)^{p} \omega_{\varphi}^{n} \tag{6}
\end{equation*}
$$

Proof. Set $T:=\omega_{\varphi}^{n-k} \wedge \omega^{k-1}$. Then

$$
\int_{X}(-\varphi)^{p-1} d \varphi \wedge d^{c} \varphi \wedge T=-\frac{1}{p} \int_{X} d\left((-\varphi)^{p}\right) \wedge d^{c} \varphi \wedge T=\frac{1}{p} \int_{X}(-\varphi)^{p} d d^{c} \varphi \wedge T
$$

Therefore the last integral is nonnegative and thus

$$
\int_{X}(-\varphi)^{p} \omega \wedge T \leq \int_{X}(-\varphi)^{p} \omega_{\varphi} \wedge T
$$

so by induction on $k$ we get (5). We also obtain

$$
\int_{X}(-\varphi)^{p-1} d \varphi \wedge d^{c} \varphi \wedge T \leq \frac{1}{p} \int_{X}(-\varphi)^{p} \omega_{\varphi} \wedge T
$$

which, by virtue of (5), gives (6).

Lemma 4. For $k=0,1, \ldots, n-1$ set $p_{k}:=1-2^{-k}$. Assume that $\varphi, \psi \in$ $\operatorname{PSH}(X, \omega) \cap L^{\infty}(X)$ are $\leq-1$ and denote $\rho:=\psi-\varphi, h:=(\varphi+\psi) / 2$. Then for $p \geq p_{n-1}$

$$
\int_{X}(-h)^{p-p_{k}} d \rho \wedge d^{c} \rho \wedge \omega_{h}^{n-1-k} \wedge \omega^{k} \leq C\left(\int_{X}(-h)^{p} d \rho \wedge d^{c} \rho \wedge \omega_{h}^{n-1}\right)^{2^{-k}}
$$

where $C$ is a positive constant depending only on $n$ and on upper bounds for $\int_{X}(-h)^{p} \omega_{h}^{n}$ and $\int_{X}(-h)^{p} d \rho \wedge d^{c} \rho \wedge \omega_{h}^{n-1}$.
Proof. We use induction on $k$. For $k=0$ there is nothing to prove and we assume the estimate holds for $k-1$. We may write the left-hand side as

$$
\int_{X}(-h)^{p-p_{k}} d \rho \wedge d^{c} \rho \wedge \omega_{h} \wedge T-\int_{X}(-h)^{p-p_{k}} d \rho \wedge d^{c} \rho \wedge d d^{c} h \wedge T
$$

where $T=\omega_{h}^{n-1-k} \wedge \omega^{k-1}$. The first integral is now estimated by the inductive assumption (and since $h \leq-1$ ), so it is enough to bound the second term from above. Note that for $q \geq 0$ we have

$$
-(-h)^{q} d d^{c} h=\frac{1}{q+1} d d^{c}\left((-h)^{q+1}\right)-q(-h)^{q-1} d h \wedge d^{c} h \leq \frac{1}{q+1} d d^{c}\left((-h)^{q+1}\right)
$$

Therefore

$$
\begin{aligned}
-\int_{X}(-h)^{p-p_{k}} d \rho \wedge d^{c} \rho & \wedge d d^{c} h \wedge T \\
& \leq \frac{1}{p-p_{k}+1} \int_{X} d \rho \wedge d^{c} \rho \wedge d d^{c}\left((-h)^{p-p_{k}+1}\right) \wedge T \\
& =-\frac{1}{p-p_{k}+1} \int_{X} d\left((-h)^{p-p_{k}+1}\right) \wedge d^{c} \rho \wedge d d^{c} \rho \wedge T \\
& =\int_{X}(-h)^{p-p_{k}} d h \wedge d^{c} \rho \wedge\left(\omega_{\psi}-\omega_{\varphi}\right) \wedge T
\end{aligned}
$$

Since $\omega_{\psi} \leq 2 \omega_{h}$, by the Schwarz inequality we get

$$
\begin{aligned}
& \left|\int_{X}(-h)^{p-p_{k}} d h \wedge d^{c} \rho \wedge \omega_{\psi} \wedge T\right| \\
& \quad \leq 2 \sqrt{\int_{X}(-h)^{p-1} d h \wedge d^{c} h \wedge \omega_{h} \wedge T} \sqrt{\int_{X}(-h)^{p-p_{k-1}} d \rho \wedge d^{c} \rho \wedge \omega_{h} \wedge T}
\end{aligned}
$$

Similarly we can deal with the term involving $\omega_{\varphi}$ and the required estimate follows from Lemma 3.

Lemma 4 easily gives Theorem 1 :
Proof of Theorem 1 for arbitrary $n$. Using the same notation as previously we can similarly as in the proof of (3) show that

$$
\lim _{j \rightarrow \infty} \int_{X}\left(-h_{j}\right)^{1-2^{1-n}} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega_{h_{j}}^{n-1}=0
$$

Lemma 4 applied for $k=n-1$, together with Lemma 2.3 in [9], now give

$$
\lim _{j \rightarrow \infty} \int_{X} d \rho_{j} \wedge d^{c} \rho_{j} \wedge \omega^{n-1}=0
$$

For the proof of Theorem 2 we will need the following quantitative version of Example 2.14 in [9]:
Proposition 5. Assume that $\psi \in \operatorname{PSH}(X, \omega)$ is negative and $0<\alpha<1$. Then $-(-\psi)^{\alpha} \in \mathcal{E}^{p}(X, \omega)$ for $p<(1-\alpha) / \alpha$.
Proof. Without loss of generality we may assume that $\psi \leq-1$. Set $\varphi:=-(-\psi)^{\alpha}$. We have

$$
\begin{aligned}
\omega_{\varphi} & =\alpha(1-\alpha)|\psi|^{\alpha-2} d \psi \wedge d^{c} \psi+\alpha|\psi|^{\alpha-1} \omega_{\psi}+\left(1-\alpha|\psi|^{\alpha-1}\right) \omega \\
& \leq \alpha(1-\alpha)|\psi|^{\alpha-2} d \psi \wedge d^{c} \psi+\alpha|\psi|^{\alpha-1} \omega_{\psi}+\omega
\end{aligned}
$$

and (for bounded $\psi$ )

$$
\begin{aligned}
\omega_{\varphi}^{n} / C \leq & \sum_{k=0}^{n-1}|\psi|^{\alpha-2+k(\alpha-1)} d \psi \wedge d^{c} \psi \wedge \omega_{\psi}^{k} \wedge \omega^{n-1-k} \\
& +\sum_{l=1}^{n}|\psi|^{l(\alpha-1)} \omega_{\psi}^{l} \wedge \omega^{n-l}+\omega^{n}
\end{aligned}
$$

where $C$ is a positive constant depending on $\alpha$ and $n$. For $a>0$ and $T=\omega_{\psi}^{k} \wedge$ $\omega^{n-1-k}$ we get

$$
\begin{aligned}
\int_{X}(-\psi)^{-a-1} d \psi \wedge d^{c} \psi \wedge T & =\frac{1}{a} \int_{X} d(-\psi)^{-a} \wedge d^{c} \psi \wedge T \\
& =-\frac{1}{a} \int_{X}(-\psi)^{-a} d d^{c} \psi \wedge T \\
& \leq \frac{1}{a} \int_{X}(-\psi)^{-a} \omega \wedge T \\
& \leq \frac{1}{a} \int_{X} \omega \wedge T \\
& =\frac{1}{a} \int_{X} \omega^{n}
\end{aligned}
$$

Therefore for $b<1-\alpha$ we obtain

$$
\int_{X}|\psi|^{b} \omega_{\varphi}^{n} \leq C(n, \alpha, b)\left(1+\int_{X}|\psi|^{b} \omega^{n}\right)
$$

and approximating arbitrary $\psi$ by $\max \{\psi,-j\}$ the proposition follows.
Proof of Theorem 2. For $z \in \mathbb{C}^{n}$ set

$$
\psi(z):=\log \left|z_{1}\right|-\frac{1}{2} \log \left(1+|z|^{2}\right)-1
$$

It can be extended to a function from $\operatorname{PSH}(X, \omega)$. Then $\widetilde{\psi}:=-\sqrt{-\psi} \notin W^{1,2}(X)$ and by Proposition $5 \widetilde{\psi} \in \mathcal{E}^{p}(X, \omega)$ for $p<1$. Using Theorem 1 we then conclude that for any $\varphi \in \mathcal{E}(X, \omega)$ satisfying $\omega_{\varphi}^{n}=\omega_{\widetilde{\psi}}^{n}$ we have $\varphi=\widetilde{\psi}+$ const $\notin W^{1,2}(X)$.

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